## ON FLOWS OF AN IDEAL GAS WHOSE SONIC SURFACE COINCIDES WITH A CHARACTERISTIC SURFACE

## (O TECHENIIAKH IDEAL'HOGO GAEA SO ZVUKOVOI POVERKHNOST'IU, SOVPADAIUSHCHEI S KHARAKTERISTICHESKOI)

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We investigate the flow of an ideal gas whose sonic surface coincides with a characteristic one (if the flow exists). For brevity, we shall call it an A-flow. In [1] a general condition for A-flows was found, namely, the sonic surface is a minimal one. In the present paper, a supplementary, necessary condition for potential A-flows is presented; in particular, it is not satisfied in the axisymmetric case (\*).

The theorem of [1] is deduced from the Bernoulli equations of continuity, of the state  $p = p(\rho, S)$  at the condition of isentropic flow. If  $n_i$  is the unit velocity vector, then

div 
$$(\rho v \mathbf{n}_1) = \frac{\partial \rho v}{\partial s_1} + \rho v \operatorname{div} \mathbf{n}_1 = 0, \qquad \frac{\partial}{\partial s_1} = \mathbf{n}_1 \cdot \nabla$$
 (1)

Since  $\partial(pv)/\partial a_1 = 0$  for v = a, it follows from (1) that div  $a_1 = 0$  for v = a; this is the condition that the surface orthogonal to  $a_1$  is a minimal one.

Now let us investigate the equation of motion  $(V \cdot \nabla) V + \rho^{-1} \nabla p = 0$ , written in the form [2]

$$v \frac{\partial v}{\partial s_1} + \frac{1}{\rho} \frac{\partial p}{\partial s_1} = 0, \quad -v^2 \varkappa + \frac{1}{\rho} \frac{\partial p}{\partial s_2} = 0, \quad \frac{\partial p}{\partial s_3} = 0 \qquad \left(\frac{\partial}{\partial s_i} = \mathbf{n}_i \cdot \nabla\right) \quad (2)$$

Here  $n_s$ ,  $n_s$  are unit vectors of the principal normal and the binormal to a streamline,  $\chi \ge 0$  is the streamline curvature.

From the second equation of (2) it follows that, if there exists a potential A-flow, then streamlines in the neighborhood of the sonic surface are approximated by straight lines, to accuracy up to third-order small quantities.

We shall determine the class of three-dimensional flows for which this necessary condition can occur.

Introduce into the potential A-flow an oblique system of coordinates so that two families,  $u_s$ ,  $u_s$ , are stream surfaces while the  $u_1$  is orthogonal to streamlines. (The surfaces  $u_1$  are assumed to be sufficiently smooth.)

\*) The fact that A-flows cannot exist in the axisymmetric case was pointed out to the author by Iu.D.Shmyglevskii (in searching for a solution of the Cauchy problem in the neighborhood of the sonic line in the form of a power series, the coefficients turned out to be imaginary). Consider a twice continuously differentiable streamline  $r(u_1)$  which intersects the sonic surface. Construct a tetrahedral elementary stream tube with constant cross-section area  $\epsilon$  so that two sides of the tube are defined by the surfaces  $u_2$  and  $u_3$ , intersecting on the curve  $r(u_1)$ , and a third side by the surface  $u_3 + 4u_3$ . Denote by

$$\frac{1}{H_3} = \frac{du_3}{|\nabla u_3|}$$

the distance on the normal between the surfaces  $u_3$  and  $u_3 + Au_3$  along the curve  $\mathbf{r}(u_1)$ ; let t be a unit vector defined on the curve  $\mathbf{r}(u_1)$  and obtained by rotating the vector  $\mathbf{a}_1$  through the angle  $\frac{1}{2}\pi$  in the given direction in the plane which is tangent to  $u_3$ .

The equation of the curve  $\mathbf{R}(u_1)$ , which lies on the surface  $u_3$  and is the edge of the elementary tube of constant cross section, may be written, with accuracy to quantities of order  $e^3$ , in the form

$$\mathbf{R}(u_1) = \mathbf{r}(u_1) + \varepsilon H_3(u_1) \mathbf{t}(u_1)$$

The area of the elementary stream tube has a minimum at the sonic point. Therefore, the streamline  $r(u_1)$  and the streamline which passes through the same point on the sonic surface as does the curve  $R(u_1)$  lie on different sides of  $R(u_1)$ ; the curvature of every streamline is zero at the sonic point; consequently, it is necessary that, at the sonic point, the projection of the curve  $R(u_1)$  on the tangent plane to the surface  $u_3$  should not be convex to the streamline  $r(u_1)$  for sufficiently small  $\epsilon$ .

If the unit vector of the curve  $\mathbf{R}(u_1)$  be denoted by  $\mathbf{N}_1$ , then this condition may be written as follows:

$$\frac{\partial \mathbf{N}_1}{\partial u_1} \cdot \mathbf{t} \leqslant 0 \qquad \text{for } v = a \tag{3}$$

Denoting a derivative with respect to  $u_1$  by a prime, we obtain

$$N_{1} = \frac{\mathbf{R}'}{|\mathbf{R}'|} = \frac{|\mathbf{r}'| \mathbf{n}_{1} + \varepsilon H_{3} \mathbf{t} + \varepsilon H_{3} \mathbf{t}'}{(|\mathbf{r}'|^{2} + \varepsilon^{2} H_{3}'^{2} + \varepsilon^{2} H_{3}^{2} \mathbf{t}' \cdot \mathbf{t}' + 2\varepsilon |\mathbf{r}'| H_{3} \mathbf{n}_{1} \cdot \mathbf{t})^{1/2}}$$

$$N_{1}' = \frac{1}{|\mathbf{R}'|} \{|\mathbf{r}'|' \mathbf{n}_{1} - |\mathbf{r}'|^{2} \times \mathbf{n}_{2} + H_{3}'' \varepsilon \mathbf{t} + H_{3} \varepsilon \mathbf{t}'' + 2\varepsilon H_{3}' \mathbf{t}' - \frac{1}{|\mathbf{R}'|^{2}} \left[|\mathbf{r}'|| \mathbf{r}'|' + \frac{\varepsilon}{2} (\varepsilon H_{3}'^{2} + \varepsilon H_{3}^{2} \mathbf{t}' \cdot \mathbf{t}' + 2 |\mathbf{r}'| H_{3} \mathbf{n}_{1} \cdot \mathbf{t}')'\right] (|\mathbf{r}'| \mathbf{n}_{1} + H_{3}' \varepsilon \mathbf{t} + H_{3} \varepsilon \mathbf{t}' + H_{3} \varepsilon \mathbf{t}'' + \frac{1}{2} \varepsilon \mathbf{t}'' + \frac{\varepsilon}{2} (\varepsilon H_{3}'' + \varepsilon H_{3}^{2} \mathbf{t}' \cdot \mathbf{t}' + 2 |\mathbf{r}'| H_{3} \mathbf{n}_{1} \cdot \mathbf{t}')'$$

We choose the family  $u_1$  so that  $|\mathbf{r}'|'=0$  on the line  $\mathbf{r}(u_1)$  at the sonic point (e.g. we put  $u_1 = s_1$ , where  $s_1$  is the arc length along the curve  $\mathbf{r}(u_1)$ ).

Let  $\gamma$  be the geodesic curvature,  $\delta$  the relative twist of the curve  $r(u_1)$  on the surface  $u_3$ ; then [3] we have

$$\mathbf{t}'' \cdot \mathbf{t} = -\mathbf{t}' \cdot \mathbf{t}' = -|\mathbf{r}'|^2 (\gamma^2 + \delta^2)$$

and condition (3) may be transformed to the form (quantities of order  $\varepsilon^2$  are neglected)

$$\frac{1}{|\nabla u_3|} \frac{\partial^2}{\partial s_1^2} |\nabla u_3| \leqslant \delta^2 \quad \text{for } v = a \tag{4}$$

In a potential A-flow, this condition holds for arbitrary choice of the family  $u_3$ .

If the flow is such that the coordinate system  $u_1$  may be chosen to be thrice orthogonal (such a system is unique, with accuracy to the order indicated, if the flow is not uniform and rectilinear), then condition (4) for the surfaces  $u_2$  and  $u_3$  can be written in the form

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$$\frac{\partial^{a}}{\partial s_{1}^{2}} |\nabla u_{i}| \leqslant 0 \quad \text{for } v = a \tag{5}$$

In the case of axial symmetry, this condition is not satisfied, as long as the sonic surface is not a plane perpendicular to the axis. In fact, we choose some streamline y = y(x) and determine  $u_1 = x$  on it. With this,

$$|\mathbf{r}'|' = \frac{y'y''}{(1+y'^2)^{1/2}} = 0$$
 for  $v = a$ ,  $|\nabla u_3| = \frac{1}{y}$ 

and condition (5) will be written in the form

 $yy'' \geqslant 2y'^2$  for v = a

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